



HAL
open science

Associated molecular liquids at the graphene monolayer interface

Florent Goujon, Aziz Ghoufi, Patrice Malfreyt

► **To cite this version:**

Florent Goujon, Aziz Ghoufi, Patrice Malfreyt. Associated molecular liquids at the graphene monolayer interface. The Journal of Chemical Physics, 2021, 154 (10), 10.1063/5.0042438 . hal-03166486

HAL Id: hal-03166486

<https://uca.hal.science/hal-03166486>

Submitted on 11 Mar 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Associated molecular liquids at the graphene monolayer interface

Florent Goujon,¹ Aziz Ghoufi,² and Patrice Malfreyt^{1, a)}¹⁾ Université Clermont Auvergne, CNRS, SIGMA Clermont, Institut de Chimie de Clermont-Ferrand (ICCF), F-63000 Clermont-Ferrand, France²⁾ Université de Rennes, CNRS, IPR (Institut de Physique de Rennes) - UMR 6251, F-35000 Rennes, France

(Dated: 2 February 2021)

We report molecular simulations of the interaction between a graphene sheet and different liquids such as water, ethanol and ethylene glycol. We describe the structural arrangements at the graphene interface in terms of density profiles, number of hydrogen bonds (HBs) and local structuration in neighboring layers close to the surface. We establish the formation of a two-dimensional HB network in the layer closest to the graphene. We also calculate the interfacial tension of liquids with a graphene monolayer and its profile along the direction normal to the graphene to rationalize and quantify the strengthening of the intermolecular interactions in the liquid due to the presence of the surface.

I. INTRODUCTION

Since graphene is used with many organic molecules in electronic organic devices¹⁻⁵, nanopore sequencing⁶, selective filtration⁷, water desalination^{8,9}, energy storage¹⁰, coating materials¹¹, strain sensors¹² and ex-foliation processes⁵, the understanding of the properties at the graphene-liquid interface is critical in terms of interactions and local arrangements. The knowledge of the local properties at the interface is an essential step toward the development of graphene-based organic devices. Indeed, the description of the first layers of molecules close to the interface are relevant for many phenomena such as charge transport processes⁴ and adsorption¹³. The structuration of the molecules at the interface is then governed by a balance between specific interactions that can be described through van der Waals, electrostatic or hydrogen bonds energy. These different energy contributions and the organization at the interface should be reflected in a key-property : the solid-liquid interfacial tension.

Indeed, the solid-liquid interfacial tension describes the balance of interactions occurring at the interfacial region. In addition, the knowledge of the liquid-vapor and solid-vapor surface tensions then allows the calculation of the contact angle θ through the Young equation¹⁴⁻¹⁶. In the case of the interaction between water and graphene, we can address the wettability of graphene (hydrophobic or hydrophilic) through the measurement of the contact angle of a water droplet. However, the measurements of the contact angle are impacted by the nature of the solid surface (roughness and/or chemical heterogeneities, liquid penetration, surface deformation) and the experimental conditions (contamination of the solid substrate by adsorption of hydrocarbons). These experimental difficulties led to scattered values¹⁷⁻²¹ of contact angle ranging from 42 to 125° giving either a hydrophobic^{18,22-24} or hydrophilic^{19,20,25} property to the graphene.

Theoretical studies also yield controversial conclusions

and do not allow to remove the ambiguity on the wetting properties of graphene. Indeed, some quantum calculations^{26,27} have predicted a water contact angle of about 85° whereas the simulated contact angle with recent atomistic simulations²⁸⁻³¹ is within the range of 56-100°.

By considering a pristine graphene sheet, it is now possible to calculate directly the graphene-water interfacial tension by molecular simulations³¹ through the stress profile. Indeed, the methodological parameters³²⁻³⁵ that have impacted the calculation of the surface tension of liquid-vapor interfaces for a long time (at least 40 years) are now under control and are in the process of being controlled for the calculation of the interfacial tension of solid-liquid interfaces^{31,36,37}.

Recently, the $\gamma_{\text{graphene-vapor}}$ interfacial tension was measured experimentally¹⁹ by using surface force apparatus. It was the first direct measurement of graphene surface energy and $\gamma_{\text{graphene-vapor}}$ found to be equal to 115 mN m⁻¹. Knowledge of this value is a key-element in the use of the Young equation¹⁴⁻¹⁶ since both interfacial tension $\gamma_{\text{liquid-vapor}}$ and $\gamma_{\text{graphene-liquid}}$ are then accessible by molecular simulation. In the case of the graphene-water interaction, we have established that $\gamma_{\text{graphene-water}}$ is close to 95 mN m⁻¹ leading to a contact angle of approximately 75° which matches very well with that resulting from surface force apparatus measurements¹⁹. The graphene surface would be a little more hydrophilic than hydrophobic. In addition, recent atomistic simulations^{31,38} clearly show that water molecules form a double layer of water molecules at the graphene surface. These layers are very well-structured thanks to a strong local network of hydrogen bonds. This atomistic description establishes that the graphene surface is wettable on an atomic scale³⁸. In addition, recent *ab initio* molecular dynamics (AIMD) simulations^{39,40} have shown that the water molecules are arranged through a two-dimensional (2D) hydrogen bond (HB) network at the liquid-vapor interface^{39,41} and increasing local ordering at the silica-water interface⁴⁰.

Combination of molecular simulations and experimental vibrational sum frequency spectroscopy (VSFS) has

^{a)} Electronic mail: Patrice.Malfreyt@uca.fr

proved to be relevant for the investigation of the behavior of water molecules and ions⁴²⁻⁴⁴ at the liquid-vapor^{45,46} and solid-liquid^{47,48} interfaces. The water molecules at the interfacial region show stronger hydrogen bonds^{49,50} and different orientations with respect to bulk conditions. Additionally, molecules that interact with the solid in the first adsorption layer can be subjected to tangential pressures of the order of hundreds of MPa^{31,37} whereas negative tangential pressure of ten MPa occur in the liquid-vapor interfaces for the same liquid molecules. These interfacial tensions find their molecular origin in a surface effect leading to a compression in the direction parallel to the surface due to the tangential pressure component.

Based on recent developments in the calculation of interfacial tension and on the ability of the simulations to provide a thorough description of the interfacial region in terms of structures, specific arrangements and orientations, we propose here to investigate the adsorption of different molecules which present various polarities and are expected to provide different contact angles⁵¹ on the graphene sheet. These molecules (ethylene glycol, ethanol and water) differ by the number and strength of the hydrogen bonds leading to variations of the liquid-vapor surface tensions⁵²⁻⁵⁶ from 22 to 72 mN m⁻¹. We aim at calculating the graphene-liquid interfacial tension of these molecules. By comparison with the liquid-vapor surface tension which we will calculate here, we will discuss the role of local arrangements at the graphene surface in relation with the macroscopic interfacial tensions.

We chose the All-Atom (AA) model in order to study more precisely the structure at the interface in terms of hydrogen bonds (donor/acceptor). We are aware that this type of model can lead to surface tensions with larger deviations from experiments compared to United Atoms (UA) models that have reproduced very well this interfacial property^{57,58}. Other alternatives are possible and atomistic models that include the polarization by using classical Drude oscillator model^{59,60} have shown to impact significantly the calculation of the work of adhesion of water on graphite and also the contact angle indirectly²⁹. However, the use of such models has shown their limitations especially in the reproduction of the liquid-vapor surface tension of salt solutions⁶¹. On solid-liquid interfaces, the models have been much less tested. The aim of this paper is not to establish the best force field that reproduces the experimental liquid-vapor and graphene-liquid interfacial tensions but rather to report the simulated interfacial tensions with the mod-

els used here and to rationalize the differences between the graphene-liquid and liquid-vapor interfacial tensions in terms of changes of structuration and local hydrogen bond network. The use of polarizable models is outside the scope of this study.

The paper is organized as follows. Section II contains the description of the model and simulation box as well as the presentation of the calculation of the graphene-liquid interfacial tension. Section III presents the discussion of our results and our conclusions are given in Section IV.

II. SIMULATION METHODOLOGY

A. Model

The total configurational energy U is defined by

$$U = U_{\text{INTRA}} + U_{\text{INTER}} \quad (1)$$

where U_{INTRA} , U_{INTER} are the intramolecular and intermolecular energy contributions, respectively. The intermolecular contributions sum the repulsion-dispersion and electrostatic energies.

The intramolecular interactions including contributions from the stretching energy, bending energy, torsion energy, and nonbonded dispersion-repulsion and electrostatic interactions are modeled with the General AMBER Force Field (GAFF2)⁶² with constrained C-H and O-H bonds.

The intermolecular energy of the system is then

$$U_{\text{INTER}} = U_{\text{LJ}} + U_{\text{ELEC}} \quad (2)$$

where the dispersion-repulsion energy is represented by the truncated Lennard-Jones (LJ) potential of

$$U_{\text{LJ}} = \sum_{i=1}^{N-1} \sum_{j>i}^N \sum_{a=1}^{N_i} \sum_{b=1}^{N_j} 4\epsilon_{ab} \left[\left(\frac{\sigma_{ab}}{r_{iajb}} \right)^{12} - \left(\frac{\sigma_{ab}}{r_{iajb}} \right)^6 \right] \quad (3)$$

where r_{iajb} is the distance between atom a in molecule i and atom b in molecule j , ϵ_{ab} is the energy parameter of the interaction and σ_{ab} is the Lennard-Jones core diameter. N_i is the number of atoms in the molecule i .

The total electrostatic potential was calculated using the Damped Reaction Field (RFD) version⁶³. For a box with orthogonal axes, the electrostatic energy calculated with the RFD formalism is written as

$$U_{\text{ELEC}} = U_{\text{RFD}} = \frac{1}{8\pi\epsilon_0} \sum_i \sum_a \sum_{j \neq i} \sum_{b \neq a} q_{ia} q_{jb} \left[\left(\frac{\text{erfc}(\alpha r_{iajb})}{r_{iajb}} + \left(\frac{\text{erfc}(\alpha r_c)}{r_c^2} + \frac{2\alpha \exp(-\alpha^2 r_c^2)}{\sqrt{\pi} r_c} \right) r_{iajb} \right) - \left(\frac{\text{erfc}(\alpha r_c)}{r_c} + \left(\frac{\text{erfc}(\alpha r_c)}{r_c^2} + \frac{2\alpha \exp(-\alpha^2 r_c^2)}{\sqrt{\pi} r_c} \right) r_c \right) + \frac{B_0(r_{iajb}^2 - r_c^2)}{2r_c^3} \right] \quad (4)$$

where B_0 is defined as

$$B_0 = \frac{2(\epsilon_1 - 1)}{2(2\epsilon_1 + 1)} \quad (5)$$

ϵ_1 is the dielectric constant outside the cavity. For water⁶⁴, we take $\epsilon_1 = 78.5$. α was taken to be 0.2365 \AA^{-1} . We have shown in a previous work³¹ that the RFD methods compares very well with the Ewald summation technique on the calculation of the graphene-water interfacial tension and the structure of water at the interface.

The LJ parameters for the graphene atom⁶⁵ are $\sigma = 3.3997 \text{ \AA}$ and $\epsilon = 0.3594 \text{ kJ mol}^{-1}$. The potential parameters including values of σ , ϵ for the water model can be found in the original paper⁶⁶. The LJ parameters for the interactions between unlike sites are calculated by using the Lorentz-Berthelot combining rules. The partial charges of the molecules have been calculated from the density-functional theory (DFT)^{67,68}(B3LYP)⁶⁹⁻⁷¹ using the Gaussian16⁷² package and the CHELPG⁷³ procedure as a grid-based method. Figure 1a shows the final symmetrized charges of each atom obtained in this work.

B. Definition of the system

The graphene-liquid interface is modeled by placing the graphene surface as a two dimensional planar material made up of a monolayer of carbon atoms at the center of the box ($z = 0$). The carbon atoms are placed in a hexagonal honeycomb lattice. The system is formed by a graphene sheet surrounded by two reservoirs of molecules. The graphene sheet was kept rigid with no intramolecular interactions between the graphene atoms. The distance $C - C$ was fixed to 1.418 \AA . The graphene sheet is formed by 2508 carbon atoms leading to a surface area $A = L_x L_y$ where $L_x = 81.05 \text{ \AA}$ and $L_y = 80.83 \text{ \AA}$.

A liquid-vapor interface has also been simulated for the same three compounds. A bulk liquid phase is first equilibrated in the constant-NVT ensemble; the box size is then increased in the z direction normal to the interface to create a gas phase. Table I summarizes the number of molecules and the box size for the graphene-liquid (SL) interface and the liquid-vapor (LV) interface.

TABLE I. Summary of the box sizes and molecule numbers for the solid-liquid (SL) and liquid-vapor (LV) system simulated in this work.

System	N_{mol}	L_x	L_y	L_z
SL water	21400	81.05	80.83	101.5
SL ethanol	6600	81.05	80.83	106.0
SL ethylene glycol	7000	81.05	80.83	106.0
LV water	1807	30.1	30.1	250
LV ethanol	4400	60.0	60.0	300.0
LV ethylene glycol	4500	60.0	60.0	300.0

The L_z dimension of the SL systems was selected from a previous study³¹ in order to avoid any dependence of the interfacial tension on the surface area. Figures 1b and 1c show two typical configurations for the SL ethylene glycol and the LV ethanol systems, respectively.

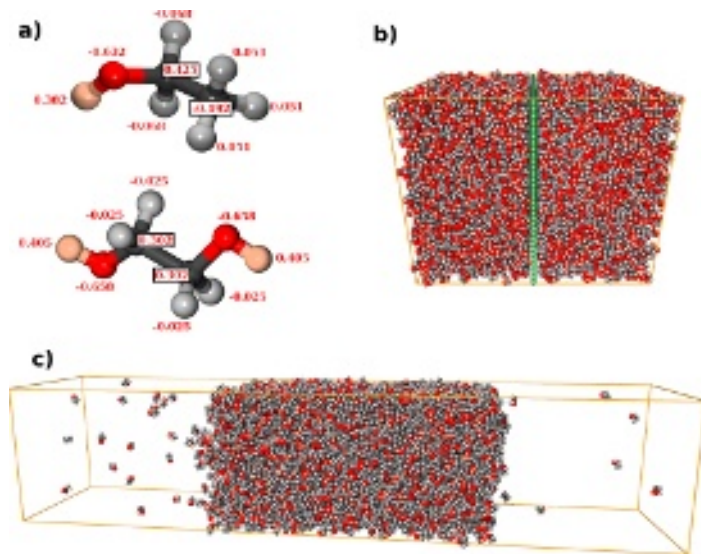


FIG. 1. a) Electrostatic charges calculated on the ethanol and ethylene glycol atoms. b) Typical configuration of a graphene layer in interaction with two reservoirs of ethylene glycol molecules. c) Typical configuration of a liquid-vapor interface of ethanol.

Molecular dynamics (MD) simulations were carried out with the DL_POLY package⁷⁴ by using the Velocity-

Verlet scheme⁷⁵ to integrate the equations of motion. Periodic boundary conditions were applied along the three directions of space. The SL simulations were carried out in the constant- Np_N AT statistical ensemble where the normal component of the pressure $p_N = p_{zz}$ was fixed to 1 bar. The area of the graphene surface is represented by A . The temperature T was fixed to 300 K. We applied the Hoover⁷⁶ thermostat and barostat with relaxation times of $\tau_t = 0.5$ ps and $\tau_p = 0.1$ ps. The LV simulations were carried out in the constant-NVT ensemble using the Hoover thermostat at $T = 300$ K with a relaxation time $\tau_t = 0.5$ ps. A time step of 2 fs is used throughout the simulations. The systems were first equilibrated over a period of 5 ns. The thermodynamic and structural properties were then averaged over an acquisition phase of 10 ns using 20000 stored configurations. The Lennard-Jones potential and the electrostatic interactions were truncated at $r_c = 12$ Å. Statistical fluctuations of interfacial tensions were estimated using the variation in the averages of five blocks of 4000 configurations.

C. Graphene-liquid interfacial tension

The graphene-liquid interfacial tension is calculated by using the mechanical definition of the interfacial tension. We use a local version which is the method of Irving and Kirkwood (IK)⁷⁷. Its expression is

$$\gamma_{IK} = \frac{1}{2} \int_{-L_z/2}^{L_z/2} (p_N(z) - p_T(z)) dz \quad (6)$$

where z indicates the direction normal to the surface and $p_N(z)$ and $p_T(z)$ are the normal and tangential components of the pressure tensor along this direction, respectively. The factor 1/2 is introduced to consider the two graphene-liquid interfaces of the system.

The method of Irving and Kirkwood⁷⁷ (IK) is based upon the notion of the force across a unit area. The components of the pressure^{78–80} tensor in the Irving and Kirkwood definition are expressed as

$$\begin{aligned} p_{\alpha\beta}(z) = & \langle \rho(z) \rangle k_B T \mathbf{I} \\ & + \frac{1}{A} \left\langle \sum_{i=1}^{N-1} \sum_{j>i}^N (\mathbf{r}_{ij})_\alpha (\mathbf{F}_{ij})_\beta \right. \\ & \left. \times \frac{1}{|z_{ij}|} \theta\left(\frac{z-z_i}{z_{ij}}\right) \theta\left(\frac{z_j-z}{z_{ij}}\right) \right\rangle \quad (7) \end{aligned}$$

where \mathbf{I} is the unit tensor and T is the input temperature. α and β represent x , y or z directions. $\theta(x)$ is the unit step function defined by $\theta(x) = 0$ when $x < 0$ and $\theta(x) = 1$ when $x \geq 0$. A is the surface area normal to the z axis. The simulation box is divided into N_z slabs of thickness $\delta z = 0.1$ Å. Following Irving and Kirkwood, atoms i and j provide a contribution to the pressure tensor in a given slab if the line joining their centres

crosses, starts or finishes in the slab. The contribution to a particular slab is $1/|z_{ij}|$ of the ij interaction. The normal component $p_N(z)$ is equal to $p_{zz}(z)$ whereas the tangential component is given by $\frac{1}{2}(p_{xx}(z) + p_{yy}(z))$. \mathbf{F}_{ij} in Eq.(7) is the intermolecular force between molecules i and j and is expressed as the sum of all the site-site forces acting between these two molecules.

The long-range corrections⁸¹ to the surface tensions γ_{LRC} are calculated using Eq.(8) by summing up the local contributions $\gamma_{LRC}(z)$ of each bin k and dividing the result by a factor 2.

$$\begin{aligned} \gamma_{LRC}(z) = & \frac{V_s}{A} \left(p_{N,lrc}(z) - p_{T,lrc}(z) \right) \\ = & \frac{\pi}{2} \rho(z) \frac{V_s}{A} \int_{r_c}^{\infty} dr \int_{-r}^r d\Delta z \frac{du_{LJ}(r)}{dr} \quad (8) \\ & \times \sum_{i=1}^{N_s} [\rho(z_i) - \rho(z_{i-1})] [r^2 - 3(\Delta z)^2] \end{aligned}$$

In the case of the graphene-water interfaces, these long-range corrections were found to be negligible³¹. For the liquid-vapor interfaces, the values of these tail corrections are given in Table II.

III. RESULTS AND DISCUSSIONS

A. Structure and orientation

The molecular density profiles along the direction z normal to the graphene sheet is given by

$$\rho(z) = \frac{1}{A\Delta z} \left\langle \sum_{i=1}^N m_i H(z_i - z, \Delta z) \right\rangle \quad (9)$$

where m_i is and z_i are the mass and z -position of the center-of-mass of molecule i , respectively. H is the top-hat function defined as

$$H(z_i - z, \Delta z) = \begin{cases} 1 & \text{if } z - \frac{\Delta z}{2} < z_i < z + \frac{\Delta z}{2} \\ 0 & \text{otherwise.} \end{cases} \quad (10)$$

A is the interfacial area, Δz is the slab thickness and z corresponds to the centers of the slabs. Δz is set to 0.05 Å for the SL systems and 0.3 Å for the LV systems.

Figure 2 displays the molecular density profiles of water, ethanol and ethylene glycol in interaction with a graphene sheet up to a distance of 25 Å from the surface in the right part of the simulation box ($z > 0$). We also give for discussion in Figure 2b the density profiles of the same species calculated in a planar liquid-vapor interface by centering each one on its Gibbs dividing surface for comparison.

When the molecules interact with a graphene sheet, the density profiles exhibit a number of adsorbed layers and amplitudes that depend on the species studied. The

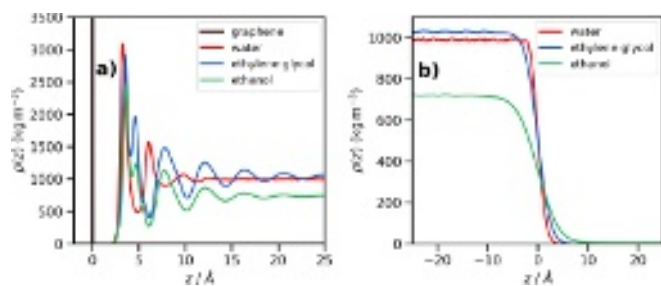


FIG. 2. Molecular density profiles of different molecules as a function of z where z represents the direction normal to the interface in the a) graphene-liquid and b) liquid-vapor interfaces.

width of the oscillations almost corresponds to the size of the molecules. The shape of the first adsorbed layer is also different between water and alcohol molecules. In the case of the liquid-vapor interfaces, the shape of the density profiles are similar between water and ethylene glycol with an interface shape close to a step-function. At the same temperature for ethanol, the interface is wider due to a smaller liquid density and an higher vapor pressure. Figure 2 shows that the molecular density profiles are radically different for species at the interface with graphene and their own vapor.

Figure 3 focuses on the density profiles of adsorbed layers on graphene by extending the region under investigation up to 40 Å from the graphene. We observe in Figure 3a a first very-well structured adsorption peak of approximately 3 times the value of the bulk liquid density. For water, the first peak is accompanied by a second peak of half the amplitude and a third slight peak located at about 10 Å from graphene. The water then recovers the bulk liquid density at a distance of 12 Å from graphene. The influence of the graphene surface is different with ethanol and ethylene glycol. Indeed, the first peak shows a shoulder whose contributions come from both carbon and oxygen atoms. There is a kind of sub-structure in this first layer for alcohol molecules due to the balance in this well-ordered layer between the strong structuration imposed by the formation of hydrogen bond (HB) and the preference for the methylene groups to interact with the graphene. The investigation of the hydrogen bond network will be proposed later in this paper. For these alcohol molecules, it is very interesting to note that the formation of structured layers extends over 20 Å and 25 Å for ethanol and ethylene glycol, respectively. This increase in the region of interaction in the z -direction is explained in part by the size of the molecule and a stronger correlation between alcohol molecules and graphene.

The layering of molecules in the interface region is also the result of an adaptation of the 3D HB network of bulk conditions to a 2D-geometry imposed by the interaction with a graphene surface. Figure 4 shows the number of hydrogen bonds per donor as function of the z -position. This number is normalized by the number of potential

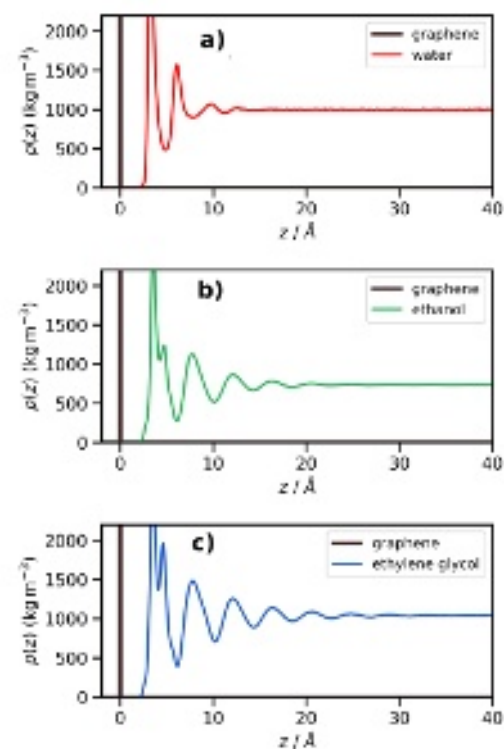


FIG. 3. Molecular density profiles $\rho(z)$ calculated over 40 Å along the z -direction for a) water, b) ethanol and c) ethylene glycol molecules. We have limited the scale on the y-axis to 2000 kg m^{-3} to have a better view of successive peaks with decreasing amplitudes.

donors (two for water and ethylene glycol and one for ethanol). The criteria used to select an hydrogen bond have been taken from Ref.82. Two water molecules are chosen as being hydrogen bonded only if the distance (O-H...O) between the donor and acceptor is less than 2.5 Å and simultaneously the angle O-H...O is greater than 135° . As each water molecule can give two HB and receive two HB, we obtain the classical value of 3.7 HB per water molecule in the bulk region by multiplying this number by 4. Interestingly, the position of the first peak of density (see Figure 3) is also the maximum of hydrogen bonds. These profiles show that the expected number of hydrogen bond by donor in bulk conditions is recovered at about 15 Å to the surface. For water close to the surface, this number is increased by 5% with respect to the bulk liquid water phase whereas an increase of 12 and 18 % is observed for ethanol and ethylene glycol with respect to bulk conditions.

Figure 5 shows the 2D-distributions of the cosine value of the angle between the O-H donor vector and the normal to the surface as a function of the position z from the graphene surface. The first layer corresponds to the first peak of density. For water, we observe a very strong orientation of water molecules in the first layer which im-

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset. PLEASE CITE THIS ARTICLE AS DOI:10.1063/1.50042438

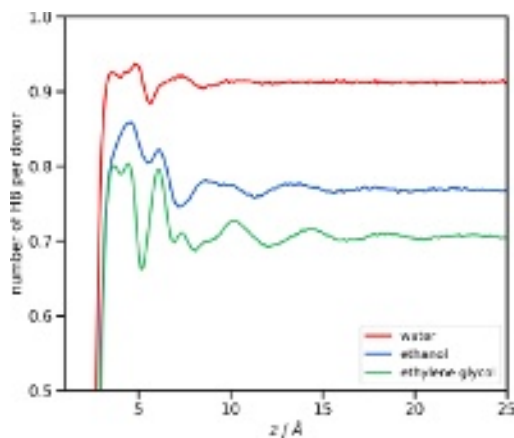


FIG. 4. Number of hydrogen bonds normalized by the number of donors (2 for water and ethylene glycol) as function of the distance to the graphene surface for water, ethylene glycol and ethanol.

plies a lesser degree of orientation in the second layer. At a distance greater than 10 Å from the graphene, the distribution of the θ angle becomes flat indicating random orientations as expected for bulk conditions. These observations are quantified in the distributions of $\cos\theta$ in Figure 6a for the first layer. The peak of the distribution for water is located at about $\cos\theta = -0.1$ suggesting that a large fraction of molecules orient the O-H vector perpendicular to the normal to the graphene. In other words, the water molecules closest to the graphene have an O-H bond parallel to the surface. This strong orientation of water molecules in the first layer also impacts the second layer by pushing the water molecules of this layer to take two preferred orientations (see Figures 5 and 6b): when the oxygen atom of water molecules of the first layer acts as an acceptor of hydrogen bond, preferred orientations occur in the second layer with O-H bond perpendicular to the graphene surface ($\cos\theta = -1$); when the O-H bond of water molecules in the first layer plays the role of donor, water molecules of the second layer are now oriented with one O-H bond both pointing toward the liquid phase with $\cos\theta = 1$. To conclude, we observe a very strong structuration of the two first layers : a first layer with hydrogen-bonds mainly oriented parallel to the surface and a second layer characterized by a majority of hydrogen bonds perpendicular to the graphene. It means that when water molecules approach the surface, there is a broken symmetry that prevents water molecules to form HBs as they do in the bulk liquid. As a result, they arrange to point the oxygen atoms toward the surface and the O-H bonds parallel to the surface forming then a 2D HB network. The number of hydrogen bonds per donor is then increased with respect to bulk conditions (see Figure 3a). Recently, many-body electronic structure methods³⁰ have shown that the orientation of a water monomer on graphene does not impact on the adsorption energy. It means that the specific orientation of water molecules

in the first layer close to the graphene monolayer is explained by the fact of maintaining a strong local network of hydrogen bonds in the interfacial region and not by the interaction with the graphene. Figure 6c shows the distributions of $\cos\theta$ in the region of the Gibbs dividing surface of a liquid-vapor interface. Of course, the distributions are wider with respect to the graphene interface but we can conclude that the water molecules exhibit a certain orientation and structuration in this interfacial region as recently shown by *ab initio* molecular dynamics calculations^{41,83,84}. The sphere of hydration is broken and the molecule reorganizes itself by creating an excess density.

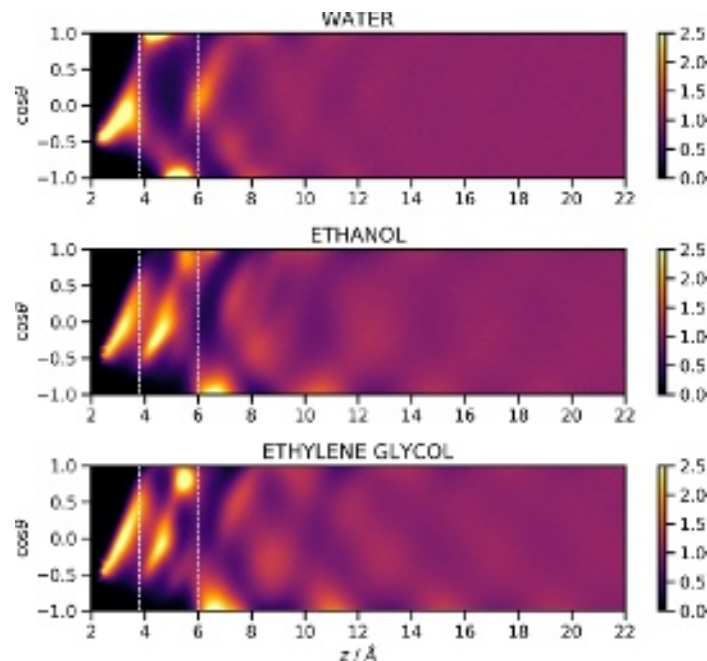


FIG. 5. 2D distribution of the OH vectors that give a hydrogen bond as a function of their position z for the three systems. Layer 1 is the closest to the graphene surface, located at $z = 0$. θ is the angle between the OH vector and the z axis normal to the surface. The colors represent the relative probability with respect to the random orientation found in the bulk region, set to unity. The dashed lines show the two first layers.

We now focus on the structuration of alcohol molecules in the graphene layer and liquid-vapor interfaces. The strong orientation observed with water molecules within a 2D organisation of hydrogen bonds parallel to the interface also applies to alcohol molecules as shown Figures 5 and 6. The main difference with water lies in a flatter angle distribution $n(\cos\theta)$ which is centered on $\langle \cos\theta \rangle = 0$. For alcohol molecules, the O-H donor has the possibility to be more parallel to the surface than with water molecules, certainly due to a less constrained network of hydrogen bonds. The second layer also shows slight differences in the distribution of θ with especially privileged orientations characterized by $\langle \cos\theta \rangle = 0$ for

ethanol and ethylene glycol and another set of favored orientations at $\langle \cos\theta \rangle = 0.75$. The 2D-HB network is maintained in the second layer, albeit to a lesser extent. This can be explained by the size of the molecules and the number of consecutive oscillating peaks in the density profiles in Figure 3. As regards the liquid-vapor interface, we also note a slight orientation of alcohol molecules parallel to the interface to a lesser degree than water. Let us recall that even in bulk conditions, alcohol molecules form chainlike structures^{85–88} which are less cohesive than the 3D-HB network of water.

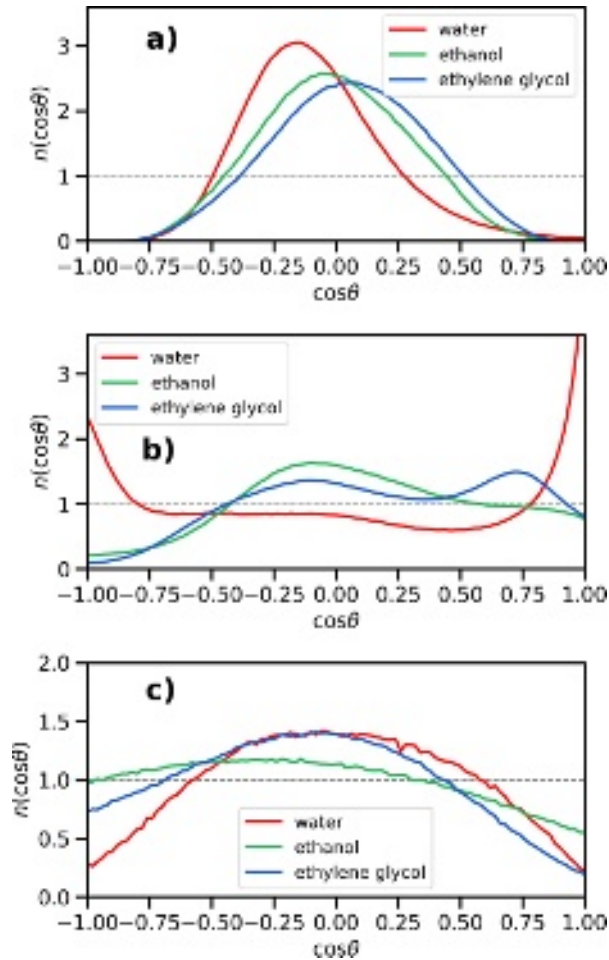


FIG. 6. $n(\cos\theta)$ distributions calculated for water and alcohol molecules for a) the first layer defined by a maximum distance to the graphene of 4 Å; b) the second layer defined by $4 < z < 6$ Å; c) a region extending over 10 Å from the position of the Gibbs dividing surface in the liquid-vapor interface.

B. Interfacial tensions

We now turn our attention on the values of interfacial tensions calculated in graphene-liquid and liquid-

TABLE II. Interfacial tensions of the liquid-vapor and graphene-water interfaces. The different contributions of the interfacial γ_{LJ} , γ_{ELEC} and γ_{LRC} are given for discussion. The average interfacial tension $\langle \gamma \rangle$ sums up the three parts. The subscripts give the accuracy of the last decimal(s), i.e., 108.7₉₀ means 108.7±9.0.

System	γ_{LJ}	γ_{ELEC}	γ_{LRC}	$\langle \gamma \rangle$ (mN m ⁻¹)
Water (SL)	-479.2	587.9	0	108.7 ₉₀
Water (LV)	-244.6	312.2	4.4	72.0 ₁₀₅
Ethanol (SL)	-24.8	92.0	0	67.2 ₁₈
Ethanol (LV)	1.5	13.3	3.9	18.7 ₁₇
Ethylene glycol (SL)	-195.0	269.3	0	74.3 ₂₆
Ethylene glycol (LV)	-105.7	137.1	6.3	37.7 ₂₈

vapor interfaces. These interfacial tensions are calculated with the Damped Reaction Field (RFD) method. The good agreement obtained with the experimental and simulated^{66,89} liquid-vapor surface tension of water at 298 K establishes that the RFD approach performs very well in the calculation of this property with the same accuracy as the Ewald method⁸⁹. In addition, the Lennard-Jones and electrostatic contributions calculated with RFD for this liquid-vapor interface are also in line with those calculated with the Ewald summation method at a higher temperature⁸⁹ ($T = 478$ K). Concerning the graphene-water interfacial tension, RFD reproduces the value calculated with Ewald³¹. For alcohol, we use the General AMBER Force Field (GAFF2)⁶² based on an all-atom description that has not been optimized on coexisting properties. This can explain some deviations from experiments on the liquid-vapor surface tensions of alcohol molecules for which United Atom models^{57,58} perform better. The aim of this paper is to rationalize the changes in the interfacial tensions between liquids and a graphene layer and liquid-vapor in line with the structural arrangements at the interface. The different contributions to the interfacial tensions are given in Table II.

From a methodological viewpoint, Figure 7c shows the profiles of the difference between the normal and tangential components of the pressure tensor and the integral $\gamma(z)$ for the water LV system. The profiles show features that are in line with the mechanical equilibrium of planar liquid-vapor interface : two peaks of about 60 MPa at the two interfaces and no contribution from the bulk liquid and vapor phases leading to constant local values of $\gamma(z)$ in these bulk phases.

The analysis is more challenging for the interaction of liquids with a graphene monolayer. Indeed, each density peak is associated with a peak of the $p_N(z) - p_T(z)$ profile. A drop of $p_N(z) - p_T(z)$ from 400 to more than -1000 MPa is calculated in the first layer of water molecules corresponding to a thickness of about 5 Å leading to a local interfacial tension of about 125 mN m⁻¹ in line with a

strong associated fluid. This profile is radically different to that observed in the liquid-vapor interface establishing that the different consecutive layers of molecules are marked by different local interfacial tensions. The strong orientation and structuration of the molecules due to the presence of the graphene clearly appear through the local values of pressures and interfacial tensions.

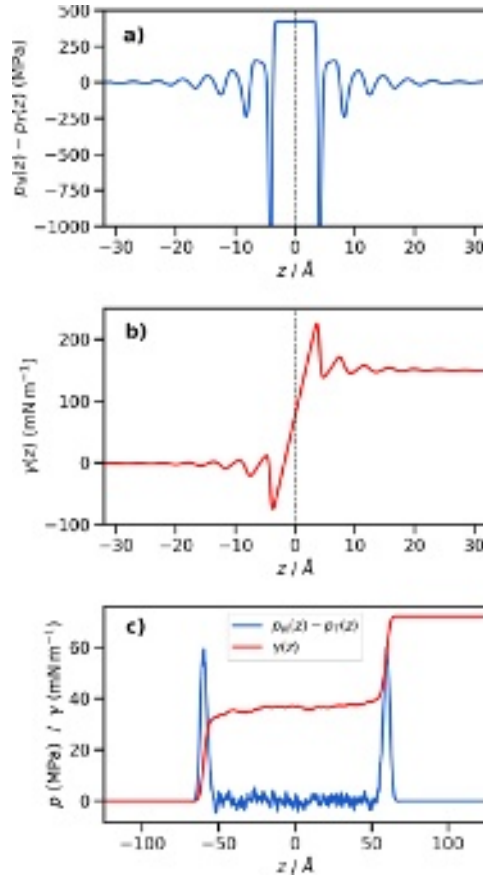


FIG. 7. Difference of $p_N(z) - p_T(z)$ (MPa) pressure profiles calculated in a) the graphene-water interface and c) the liquid-vapor interface of water. The integral $\gamma(z)$ of the local interfacial tension (mN m^{-1}) calculated using Eq.(6) is given in b) for graphene-water and in c) for water liquid-vapor systems.

For completeness, Table II lists the different contributions to the surface tensions. A negative contribution to the surface tension from the Lennard-Jones (LJ) part means that the LJ interacting centers are on average at a smaller distance than the σ diameter parameter leading to positive values for the energy. This repulsive LJ energy is then explained by strong electrostatic interactions. The magnitude of the LJ part to the interfacial tension measures to a certain extent the strength of the association in the liquid. For water and ethylene glycol, the LJ contribution to the surface tension becomes much more negative when we move from the liquid-vapor to the graphene-water interfaces. For ethanol, this contribution changes from a positive to a negative value but it is also

the lowest value of the systems studied here. From an energetic viewpoint, the strengthening of the association in the fluid due to the graphene is the most significant for water with $\gamma_{\text{SL}} = 108.7 \text{ mN m}^{-1}$ and the lightest for ethanol with $\gamma_{\text{SL}} = 67.2 \text{ mN m}^{-1}$. We do not see significant changes in the structuration of the two closest layers next to the graphene for the alcohol molecules in line with similar interfacial tensions.

IV. CONCLUSIONS

The graphene-liquid interface was investigated by means of molecular simulations. The density profiles along the direction normal to the surface establish different adsorbed layers with a number and amplitudes that depend on the molecule. The shape of the profiles changes radically from those calculated in a liquid-vapor interface.

The simulations have evidenced a very strong structuration in the two layers closest to the graphene. As water and alcohol molecules approach the graphene, they undergo a break in symmetry that cancels any possibility of forming hydrogen bond on the side of graphene. Molecules adapt to this situation by orienting the hydrogen bonds in the direction parallel to the surface leading to a 2D HB network in the first layer. The interfacial region of the liquid-vapor equilibrium also shows the formation of 2D HB network but to a lesser extent compared to graphene-liquid interface.

The layering of water and ethanol molecules when they interact with the graphene is noticeable on the profiles of interfacial tension and pressure components. In the case of water, the first layer of water sees a drop of $p_N - p_T$ of about 1500 MPa over a thickness of 5 Å. In this first layer, the local interfacial tension is increased by 74% with respect to the surface tension of liquid-vapor interface. From an energetic viewpoint, the enhancement of the interaction within the liquid due to the graphene is the most significant for water and the smallest for ethanol.

This paper aimed at showing how properties of associated liquids characterized by different hydrogen bond strengths change with respect to bulk conditions when they interact with a surface. The structural changes due to the adsorption have been analyzed in terms of local molecular density and orientations. Some dramatic changes in the pressure tensor and local surface tensions have been identified in the interfacial region with the solid. These structural and energetic modifications are significantly different from those observed in the interfacial region of the liquid-vapor equilibrium.

DATA AVAILABILITY

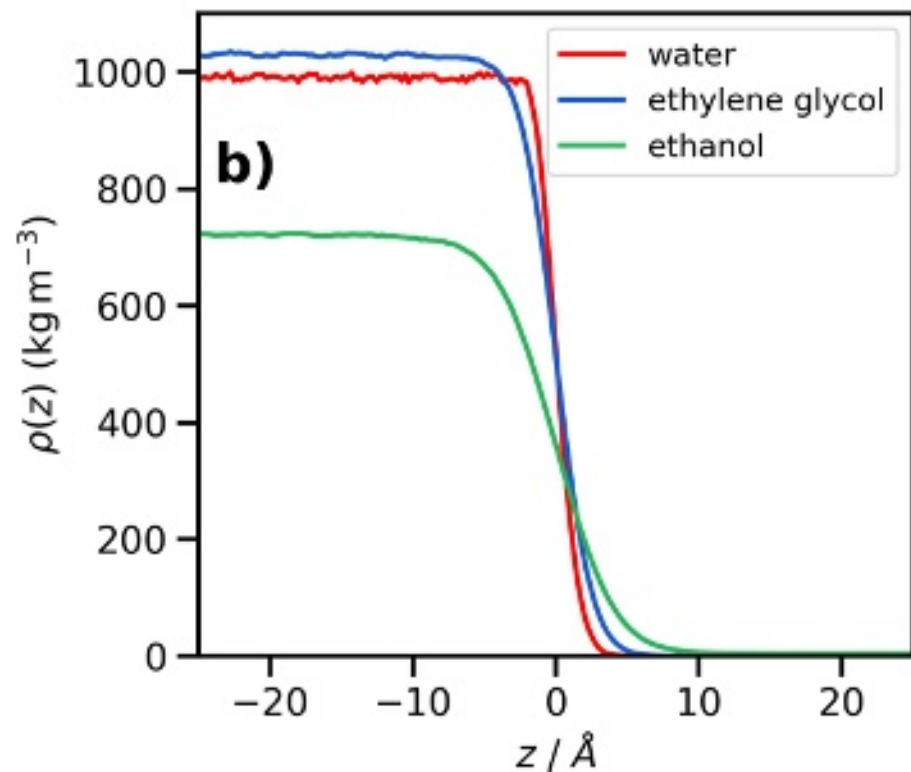
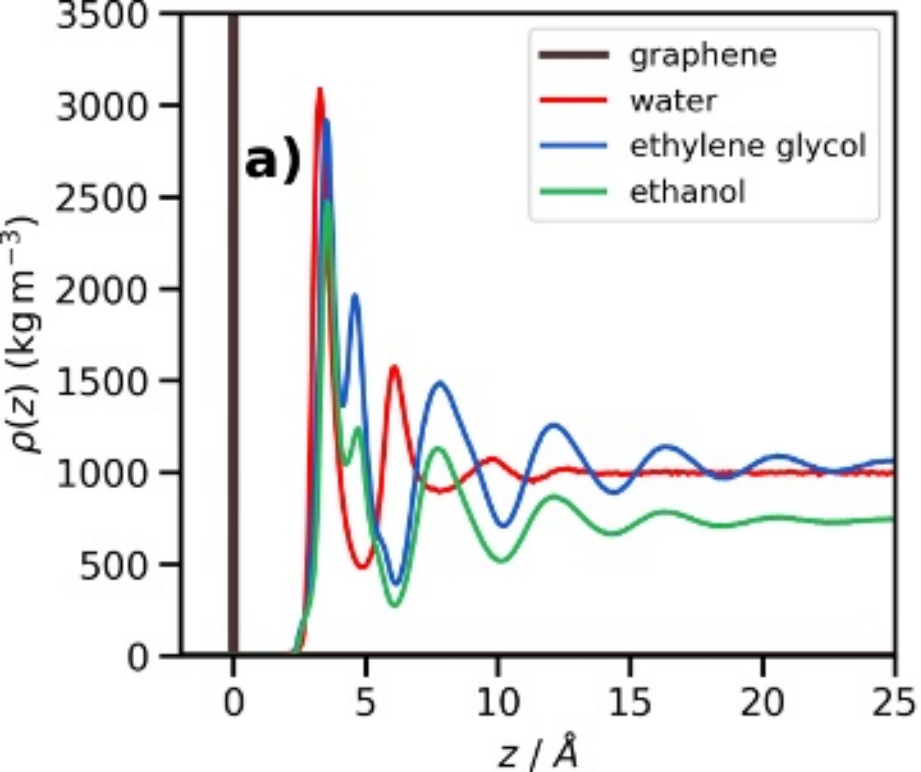
The data that support the findings of this study are available from the corresponding author upon reasonable request.

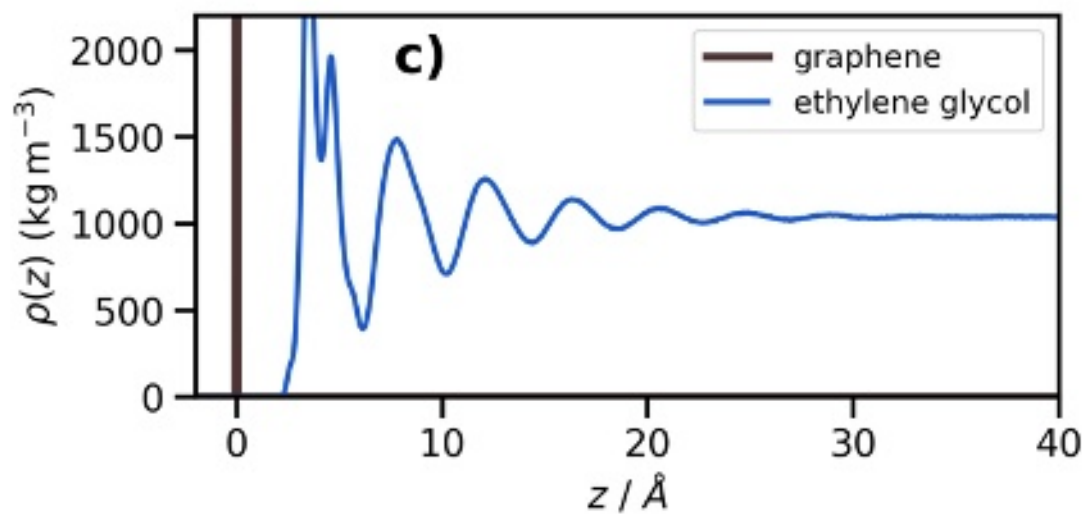
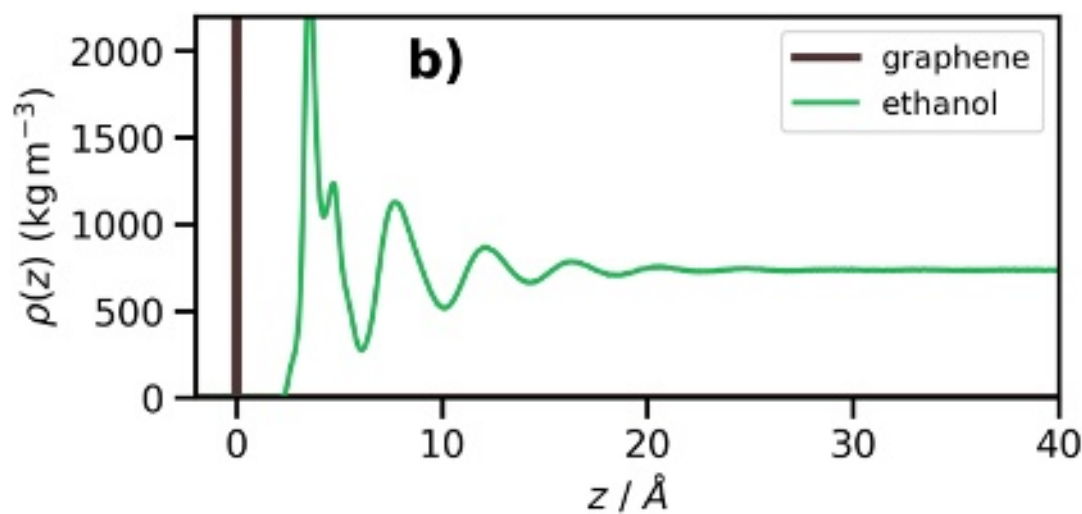
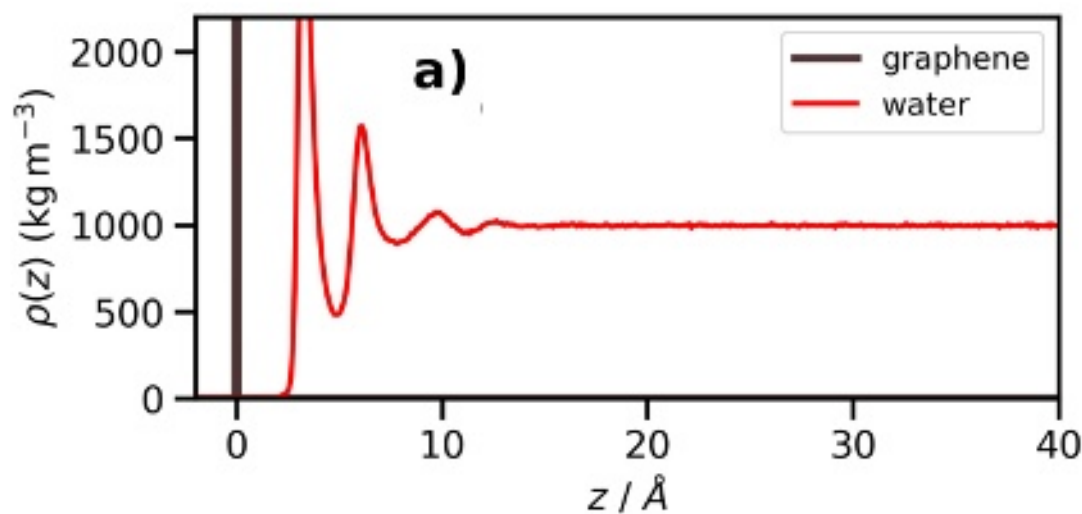
- ¹S. Coe, W.-K. Woo, M. Bawendi, and V. Bulovic, "Electroluminescence from single monolayers of nanocrystals in molecular organic devices," *Nature* **420**, 800–803 (2002).
- ²F. Bonaccorso, Z. Sun, T. Hasan, and A. C. Ferrari, "Graphene photonics and optoelectronics," *Nat. Photonics* **4**, 611–622 (2010).
- ³K. Kim, J. Y. Choi, T. Kim, S. H. Cho, and H. J. Chung, "A role for graphene in silicon-based semiconductor devices," *Nature* **479**, 338–344 (2011).
- ⁴G. Hong, Q. H. Wu, J. Ren, C. Wang, W. Zhang, and S. Lee, "Recent progress in organic molecule/graphene interfaces," *Nano Today* **8**, 388–402 (2013).
- ⁵A. Schlierf, P. Samori, and V. Palermo, "Graphene–organic composites for electronics: Optical and electronic interactions in vacuum, liquids and thin solid films," *J. Mater. Chem. C* **2** (2014).
- ⁶S. Garaj, W. Hubbard, A. Reina, J. Kong, D. Branton, and J. A. Golovchenko, "Graphene as a subnanometre trans-electrode membrane," *Nature* **467**, 190–193 (2010).
- ⁷S. C. O'Hern, C. A. Stewart, M. S. H. Boutilier, J. C. Idrobo, S. Bhaviripudi, S. K. Das, J. Kong, T. Laoui, M. Atieh, and R. Karnik, "Selective molecular transport through intrinsic defects in a single layer of cvd graphene," *ACS Nano* **6**, 10130–10138 (2012).
- ⁸D. Cohen-Tanugi and J. C. Grossman, "Water desalination across nanoporous graphene," *Nano Lett.* **12**, 3602–3608 (2012).
- ⁹L. Garnier, A. Szymczyk, P. Malfreyt, and A. Ghoufi, "Physics behind water transport through nanoporous boron nitride and graphene," *J. Phys. Chem. Lett.* **7**, 3371–3376 (2016).
- ¹⁰Y. Zhu, S. Murali, M. D. Stollernd, K. J. Ganesh, W. Cai, P. J. Ferreira, A. Pirkle, R. M. Wallace, K. A. Cychoz, M. Thomme, D. Suand, E. A. Stach, and R. S. Ruoff, "Carbon-based supercapacitors produced by activation of graphene," *Science* **332**, 1537–1541 (2011).
- ¹¹M. Topsakal, H. Sahin, and S. Ciraci, "Graphene coatings: An efficient protection from oxidation," *Phys. Rev. B* **85**, 155445 (2012).
- ¹²M. Oliva-Leyva and C. Wang, "Low-energy theory for strained graphene: an approach up to second-order in the strain tensor," *J. Phys.: Condens. Matter* **29**, 165301 (2017).
- ¹³G. Ersan, O. G. Apul, F. Perreault, and T. Karanfil, "Adsorption of organic contaminants by graphene nanosheets: A review," *Water Res.* **126**, 385–398 (2017).
- ¹⁴T. Young, "An essay on the cohesion of fluids," *Philos. Trans. R. Soc. London* **95**, 65 (1805).
- ¹⁵A. Marmur, "Soft contact: measurement and interpretation of contact angles," *Soft Matter* **2**, 12–17 (2006).
- ¹⁶E. Bormashenko, "Physics of solid–liquid interfaces: From the young equation to the superhydrophobicity," *Low Temperature Physics* **42**, 622–635 (2016).
- ¹⁷A. Ashraf, Y. Wu, M. C. Wang, N. R. Aluru, S. A. Dastgheib, and S. Nam, "Spectroscopic investigation of the wettability of multilayer graphene using highly ordered pyrolytic graphite as a model material," *Langmuir* **30**, 12827–12836 (2014).
- ¹⁸T. Ondarçuhu, V. Thomas, M. Nunez, E. Dujardin, A. Rahman, C. T. Black, and A. Checco, "Wettability of partially suspended graphene," *Sci. Rep.* **6**, 24237 (2016).
- ¹⁹C. D. van Engers, N. E. A. Cousens, V. Babenko, J. Britton, B. Zappone, N. Grobert, and S. Perkin, "Direct measurement of the surface energy of graphene," *Nano Lett.* **17**, 3815–3821 (2017).
- ²⁰A. V. Prydatko, L. A. Belyaeva, L. Jiang, L. M. Lima, and G. F. Schneider, "Contact angle measurement of free-standing square-millimeter single-layer graphene," *Nat. Commun.* **1** (2018).
- ²¹J. Feng and Z. Guo, "Wettability of graphene: from influencing factors and reversible conversions to potential applications," *Nanoscale Horiz.* **4**, 339–364 (2019).
- ²²S. Wang, Y. Zhang, N. Abidi, and L. Cabrales, "Wettability and surface free energy of graphene films," *Langmuir* **25**, 11078–11081 (2009).
- ²³G.-T. Kim, S.-J. Gim, S.-M. Cho, N. Koratkar, and I.-K. Oh, "Wetting-transparent graphene films for hydrophobic water-harvesting surfaces," *Adv. Mater.* **26**, 5166–5172 (2014).
- ²⁴M. Munz, C. E. Giusca, R. L. Myers-Ward, D. K. Gaskill, and O. Kazakova, "Thickness-dependent hydrophobicity of epitaxial graphene thickness-dependent hydrophobicity of epitaxial graphene," *ACS Nano* **9**, 8401–8411 (2015).
- ²⁵A. Kozbial, Z. Li, C. Conaway, R. McGinley, S. Dhingra, V. Vahdat, F. Zhou, B. D'Urso, H. Liu, and L. Li, "Study on the surface energy of graphene by contact angle measurements," *Langmuir* **30**, 8598–8606 (2014).
- ²⁶H. Li and X. C. Zeng, "Wetting and interfacial properties of water nanodroplets in contact with graphene and monolayer boron nitride sheets," *ACS Nano* **6**, 2401–2409 (2012).
- ²⁷J. Li and F. Wang, "Water graphene contact surface investigated by pairwise potentials from force-matching paw-pbe with dispersion correction," *J. Chem. Phys.* **146**, 054702 (2017).
- ²⁸F. Taherian, V. Marcon, N. F. A. van der Vegt, and F. Leroy, "What is the contact angle of water on graphene?" *Langmuir* **29**, 1457–1465 (2013).
- ²⁹R. P. Misra and D. Blankschtein, "Insights on the role of many-body polarization effects in the wetting of graphitic surfaces by water," *J. Phys. Chem. C* **121**, 28166–28179 (2017).
- ³⁰J. G. Brandenburg, A. Zen, M. Fitzner, B. Ramberge, G. Kresse, heodoros Tsatsoulis, A. Grüneis, A. Michaelides, and D. Alfe, "Physisorption of water on graphene: Subchemical accuracy from many-body electronic structure methods," *J. Phys. Chem. Lett.* **10**, 358–368 (2019).
- ³¹T. Dreher, C. Lemarchand, N. Pineau, E. Bourasseau, A. Ghoufi, and P. Malfreyt, "Calculation of the interfacial tension of the graphene-water interaction by molecular simulations," *J. Chem Phys.* **150**, 014703 (2019).
- ³²F. Goujon, P. Malfreyt, A. Boutin, and A. H. Fuchs, "Vapour-liquid phase equilibria of n-alkanes by direct monte carlo simulations," *Mol. Simul.* **27**, 99–114 (2001).
- ³³A. Ghoufi, F. Goujon, V. Lachet, and P. Malfreyt, "Multiple histogram reweighting method for the surface tension calculation," *J. Chem. Phys.* **128**, 154718 (2008).
- ³⁴X. Paredes, J. Fernandez, A. A. H. Padua, P. Malfreyt, F. Malberg, B. Kirchner, and A. S. Pensado, "Bulk and liquid-vapor interface of pyrrolidinium-based ionic liquids: A molecular simulation study," *J. Phys. Chem. B* **118**, 731–742 (2014).
- ³⁵A. Ghoufi, P. Malfreyt, and D. J. Tildesley, "Computer modelling of the surface tension of the gas-liquid and liquid-liquid interface," *Chem. Soc. Rev.* **45**, 1387–1409 (2016).
- ³⁶H. D. d'Oliveira, X. Davoy, E. Arche, P. Malfreyt, and A. Ghoufi, "Test-area surface tension calculation of the graphene-methane interface: Fluctuations and commensurability," *J. Chem. Phys.* **146**, 214112 (2017).
- ³⁷T. Dreher, C. Lemarchand, L. Soulard, E. Bourasseau, P. Malfreyt, and N. Pineau, "Calculation of a solid/liquid surface tension: A methodological study," *J. Chem. Phys.* **148**, 034702 (2018).
- ³⁸A. Akaishi, T. Yonemaru, and J. Nakamura, "Formation of water layers on graphene surfaces," *ACS Omega* **2**, 2184–2190 (2017).
- ³⁹A. Serva, S. Pezzotti, S. Bougueroua, D. R. Galimberti, and M.-P. Gaijeot, "Combining ab-initio and classical molecular dynamics simulations to unravel the structure of the 2d-hb-network at the air-water interface," *J. Mol. Struct.* **1165**, 71–78 (2018).
- ⁴⁰M. Gaijeot and M. Sulpizi, *Structure and Dynamics of Solid/Liquid Interfaces: Liquid and Biological Interfaces*, edited by K. Wandelt, Vol. 7 (Wiley-VCH Verlag GmbH & Co, 2020).
- ⁴¹S. Pezzotti, D. R. Galimberti, and M.-P. Gaijeot, "2d h-bond network as the topmost skin to the air-water interface," *J. Phys. Chem. Lett.* **8**, 3133–3141 (2017).
- ⁴²P. B. Petersen and R. J. Saykally, "On the nature of ions at the liquid water surface," *Annu. Rev. Phys. Chem.* **57**, 333–364 (2006).
- ⁴³D. L. McCaffrey, S. C. Nguyen, S. J. Cox, H. Weller, A. P. Alivisatos, P. L. Geissler, and R. J. Saykally, "Mechanism

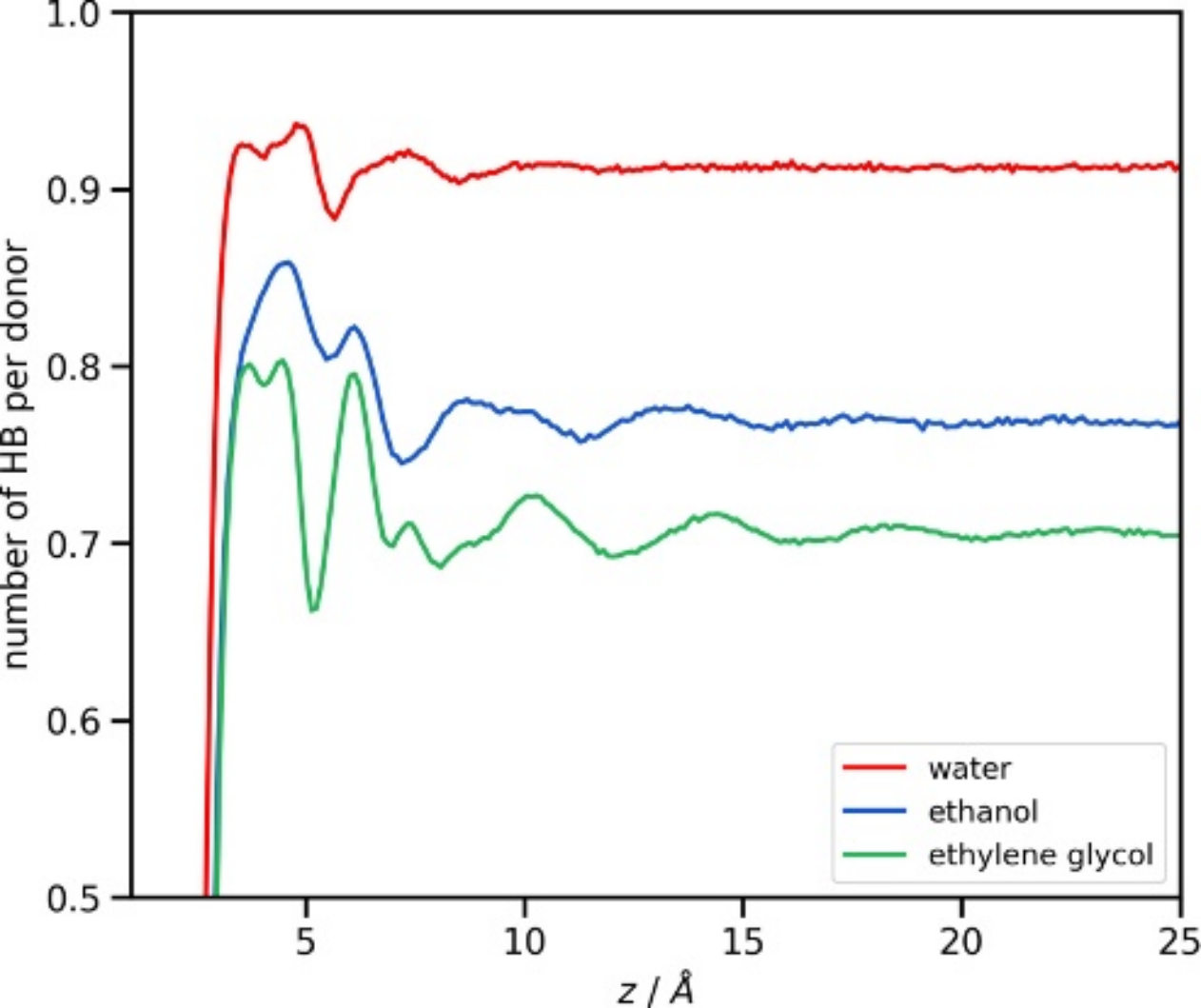
- of ion adsorption to aqueous interfaces: Graphene/water vs. air/water,” *Proc. Natl. Acad. Sci.* **114**, 13369–13373 (2017).
- ⁴⁴H. Mizuno, A. M. Rizzuto, and R. J. Saykally, “Charge-transfer-to-solvent spectrum of thiocyanate at the air/water interface measured by broadband deep ultraviolet electronic sum frequency generation spectroscopy,” *J. Phys. Chem. Lett.* **9**, 4753–4757 (2018).
- ⁴⁵E. A. Raymond, T. L. Tarbuck, M. G. Brown, and G. L. Richmond, “Hydrogen-bonding interactions at the vapor/water interface investigated by vibrational sum-frequency spectroscopy of $\text{H}_2\text{O}/\text{D}_2\text{O}$ mixtures and molecular dynamics simulations,” *J. Phys. Chem. B* **107**, 546–556 (2003).
- ⁴⁶D. S. Walker and G. L. Richmond, “Understanding the effects of hydrogen bonding at the vapor-water interface: Vibrational sum frequency spectroscopy of $\text{H}_2\text{O}/\text{D}_2\text{O}$ mixtures studied using molecular dynamics simulations,” *J. Phys. Chem. C* **111**, 8321–8330 (2007).
- ⁴⁷A. J. Hopkins, C. L. McFearn, and G. L. Richmond, “Investigations of the solid–aqueous interface with vibrational sum-frequency spectroscopy,” *Curr. Opin. Solid State Mater. Sci.* **9**, 19–27 (2005).
- ⁴⁸S. Schrödle, F. G. Moore, and G. L. Richmond, “Surface speciation at solid/liquid interfaces: A vibrational sum-frequency study of acetate adsorption at the fluorite/water interface,” *J. Phys. Chem. C* **111**, 10088–10094 (2007).
- ⁴⁹W. T. S. Cole and R. J. Saykally, “Hydrogen bond network rearrangement dynamics in water clusters: Effects of intermolecular vibrational excitation on tunneling rates,” *J. Chem. Phys.* **147**, 064301 (2017).
- ⁵⁰W. T. S. Cole, R. S. Fellers, M. R. Viant, and R. J. Saykally, “Hydrogen bond breaking dynamics in the water pentamer: Terahertz vrt spectroscopy of a 20 μm libration,” *J. Chem. Phys.* **146**, 014306 (2017).
- ⁵¹H. Wang, C. Junghans, and K. Kremer, “Comparative atomistic and coarse-grained study of water: what do we lose by coarse-graining?” *Eur. Phys. J. E* **28**, 221–229 (2009).
- ⁵²G. Vazquez, E. Alvarez, and J. M. Navaza, “Surface tension of alcohol + water from 20 to 50 $^\circ\text{C}$,” *J. Chem. Eng. Data* **40**, 611–614 (1995).
- ⁵³S. Azizian and M. Hemmati, “Surface tension of binary mixtures of ethanol + ethylene glycol from 20 to 50 $^\circ\text{C}$,” *J. Chem. Eng. Data* **48**, 662–663 (2003).
- ⁵⁴M. A. Postigo, J. Zurita, M. L. G. A. de Soria, and M. Katz, “Excess thermodynamic properties of n-pentane + dichloromethane system at 298.15 k,” *Can. J. Chem.* **64**, 1966–1968 (2011).
- ⁵⁵A. Zdziennicka, J. Krawczyk, K. Szymczyk, and B. Janczuk, “Components and parameters of liquids and some polymers surface tension at different temperature,” *Colloids and Surfaces A* **529**, 864–875 (2017).
- ⁵⁶A. Erfani, S. Khosharay, and C. P. Aichele, “Surface tension and interfacial compositions of binary glycerol/alcohol mixtures,” *J. Chem. Thermodynamics* **135**, 241–251 (2019).
- ⁵⁷F. Biscay, A. Ghoufi, and P. Malfreyt, “Surface tension of water-alcohol mixtures from monte carlo simulations,” *J. Chem. Phys.* **134**, 044709 (2011).
- ⁵⁸F. Biscay, A. Ghoufi, V. Lachet, and P. Malfreyt, “Prediction of the surface tension of the liquid-vapor interface of alcohols from monte carlo simulations,” *J. Phys. Chem. C* **115**, 8670–8683 (2011).
- ⁵⁹G. Lamoureux, A. D. MacKerell, and B. Roux, “A simple polarizable model of water based on classical drude oscillators,” *J. Chem. Phys.* **119**, 5185–5197 (2003).
- ⁶⁰G. Lamoureux, E. Harder, I. V. Vorobyov, B. Roux, and J. A. D. MacKerell, “A polarizable model of water for molecular dynamics simulations of biomolecules,” *Chem. Phys. Lett.* **418**, 245–249 (2006).
- ⁶¹J. C. Neyt, A. Wender, V. Lachet, A. Ghoufi, and P. Malfreyt, “Prediction of the concentration dependence of the surface tension and density of salt solutions: atomistic simulations using drude oscillator polarizable and nonpolarizable models,” *Phys. Chem. Chem. Phys.* **15**, 11679–11690 (2013).
- ⁶²D. Case, R. Betz, D. Cerutti, T. C. III, T. Darden, R. Duke, T. Giese, H. Gohlke, A. Goetz, N. Homeyer, S. Izadi, P. Janowski, J. Kaus, A. Kovalenko, T. Lee, S. LeGrand, P. Li, C. Lin, T. Luchko, R. Luo, B. Madej, D. Mermelstein, K. Merz, G. Monard, H. Nguyen, H. Nguyen, I. Omelyan, A. Onufriev, D. Roe, A. Roitberg, C. Sagui, C. Simmeling, W. Botello-Smith, J. Swails, R. Walker, J. Wang, R. Wolf, X. Wu, L. Xiao, and P. Kollman, “Amber 2016,” (University of California, San Francisco, 2016).
- ⁶³C. J. Fennell and J. D. Gezelter, “Is the ewald summation still necessary? pairwise alternatives to the accepted standard for long-range electrostatics,” *J. Chem. Phys.* **124**, 234104 (2006).
- ⁶⁴J. M. Míguez, M. M. Piñeiro, and F. J. Blas, “Influence of the long-range corrections on the interfacial properties of molecular models using Monte Carlo simulation,” *J. Chem. Phys.* **138**, 34707–34716 (2013).
- ⁶⁵D. Cohen-Tanugi and J. C. Grossman, “Water desalination across nanoporous graphene,” *Nano Lett.* **12**, 3602–3608 (2012).
- ⁶⁶J. L. F. Abascal and C. Vega, “A general purpose model for the condensed phases of water: TIP4P/2005,” *J. Chem. Phys.* **123**, 234505 (2005).
- ⁶⁷P. Honenberg and W. Kohn, “Inhomogeneous electron gas,” *Phys. Rev. A* **136**, 864–871 (1964).
- ⁶⁸W. Kohn and L. Sham, “Self-consistent equations including exchange and correlation effects,” *J. Phys. Rev. A* **140**, 1133–1138 (1965).
- ⁶⁹A. D. Becke, “Density-functional exchange-energy approximation with correct asymptotic behavior,” *Phys. Rev. A* **38**, 3098 (1988).
- ⁷⁰A. D. Becke, “Density-functional thermochemistry. iii. the role of exact exchange,” *J. Chem. Phys.* **98**, 5648–5652 (1993).
- ⁷¹J. P. Perdew and Y. Wang, “Accurate and simple analytic representation of the electron-gas correlation energy,” *Phys. Rev. E* **45**, 13244 (1991).
- ⁷²M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. M. Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman, and D. J. Fox, “Gaussian 16, revision b.01,” Gaussian, Inc., Wallingford CT (2016).
- ⁷³C. P. Breneman and K. B. Wiberg, “Determining atom-centered monopoles from molecular electrostatic potentials. the need for high sampling density in formamide conformational analysis,” *J. Comput. Chem.* **11**, 361–373 (1990).
- ⁷⁴I. Todorov and W. S. and, “DLPOLY3: New dimensions in molecular dynamics simulations via massive parallelism,” *J. Mater. Chem.* **16**, 1911–1918 (2006).
- ⁷⁵M. P. Allen and D. J. Tildesley, *Computer Simulation of Liquids, Second edition* (Oxford: Clarendon Press, 2017).
- ⁷⁶S. Melchionna, G. Ciccotti, and B. L. Holian, “Hoover NPT dynamics for systems varying in shape and size,” *Mol. Phys.* **78**, 533–544 (1993).
- ⁷⁷J. H. Irving and J. Kirkwood, “The Statistical Mechanical Theory of Transport Processes. IV. The Equations of Hydrodynamics,” *J. Chem. Phys.* **18**, 817–829 (1950).
- ⁷⁸J. S. Rowlinson and B. Widom, *Molecular Theory of Capillarity* (Clarendon Press, Oxford, 1982).

This is the author's peer reviewed, accepted manuscript. However, the online version of record will be different from this version once it has been copyedited and typeset.
PLEASE CITE THIS ARTICLE AS DOI:10.1063/1.50042438

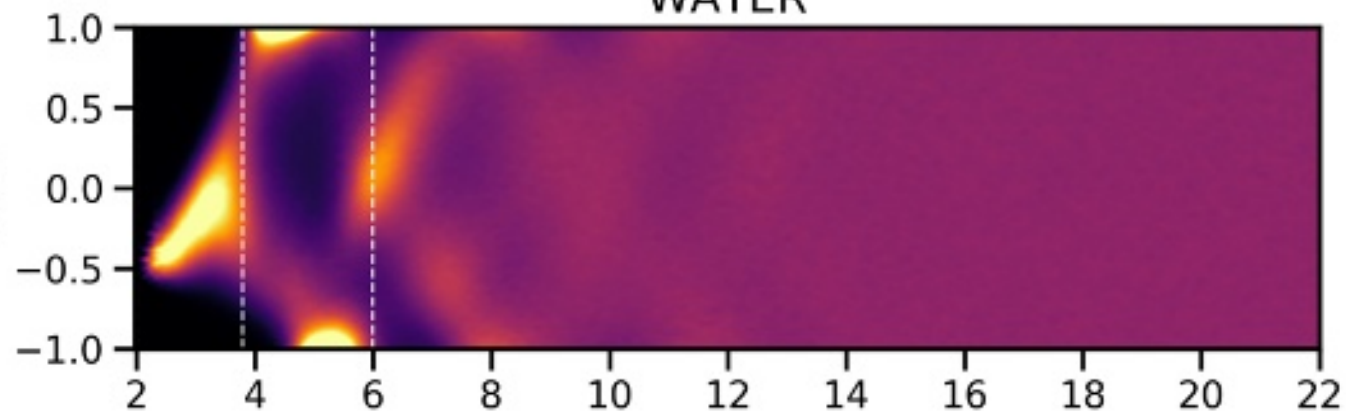
- ⁷⁹J. P. R. B. Walton, D. J. Tildesley, J. S. Rowlinson, and J. R. Henderson, "The pressure tensor at the planar surface of a liquid," *Molec. Phys.* **48**, 1357–1368 (1983).
- ⁸⁰J. P. R. B. Walton, D. J. Tildesley, J. S. Rowlinson, and J. R. Henderson, "The pressure tensor at the planar surface of a liquid," *Molec. Phys.* **50**, 1381 (1983), erratum.
- ⁸¹M. Guo and B. Lu, "Long range corrections to thermodynamic properties of inhomogeneous systems with planar interfaces," *J. Chem. Phys.* **106**, 3688–3695 (1997).
- ⁸²A. Luzar and D. Chandler, "Effect of environment on hydrogen bond dynamics in liquid water," *Phys. Rev. Lett.* **76**, 928–931 (1996).
- ⁸³A. P. Willard and D. Chandler, "Instantaneous liquid interfaces," *J. Phys. Chem. B* **114**, 1954–1958 (2010).
- ⁸⁴J. Kessler, H. Elgabarty, T. Spura, K. Karhan, P. Partovi-Azar, A. A. Hassanali, and T. D. Kühne, "Structure and dynamics of the instantaneous water/vapor interface revisited by path-integral and ab initio molecular dynamics simulations," *J. Phys. Chem. B* **119**, 10079–10086 (2015).
- ⁸⁵T. E. Goresy and R. Böhmer, "Diluting the hydrogen bonds in viscous solutions of *n*-butanol with *n*-bromobutane: A dielectric study," *J. Chem. Phys.* **128**, 154520 (2008).
- ⁸⁶R. Böhmer, C. Gainarua, and R. Richert, "Structure and dynamics of monohydroxy alcohols—milestones towards their microscopic understanding, 100 years after debye," *Phys. Rep.* **545**, 125–195 (2014).
- ⁸⁷M. L. Liriano, J. Carrasco, E. A. Lewis, C. J. Murphy, T. J. Lawton, M. D. Marcinkowski, A. J. Therrien, A. Michaelides, and E. C. H. Sykes, "The interplay of covalency, hydrogen bonding, and dispersion leads to a long range chiral network: The example of 2-butanol," *J. Chem Phys.* **144**, 094703 (2016).
- ⁸⁸I. Essafri and A. Ghoufi, "Microstructure of nonideal methanol binary liquid mixtures," *Phys. Rev. E* **99**, 062607 (2019).
- ⁸⁹A. Ghoufi, F. Goujon, V. Lachet, and P. Malfreyt, "Surface tension of water and acid gases from Monte Carlo simulations," *J. Chem. Phys.* **128**, 154716–154731 (2008).



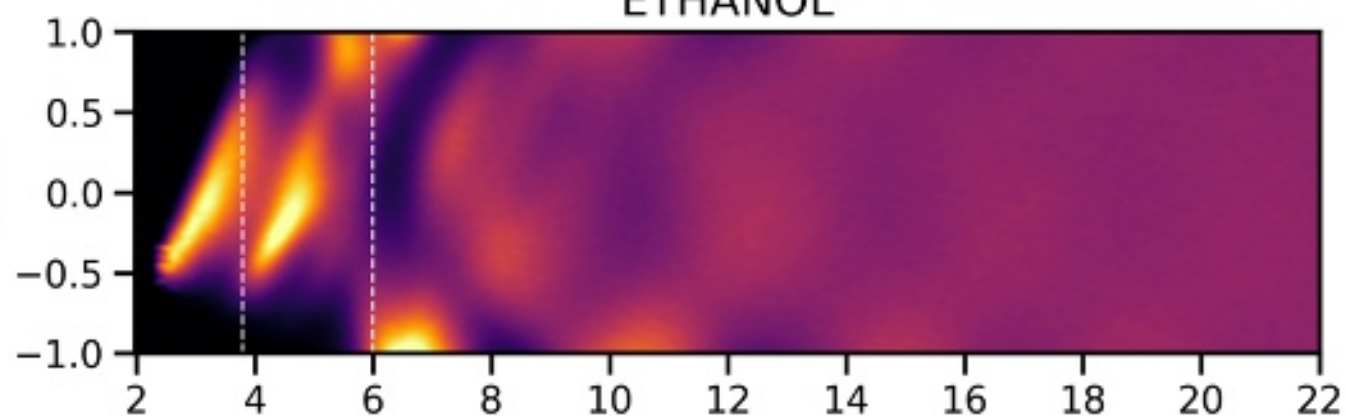




WATER



ETHANOL



ETHYLENE GLYCOL

